

On Reconstruction of the Dynamic Tortuosity Functions of Poroelastic Materials

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**SYMPOSIUM ON THE ACOUSTICS OF PORO-ELASTIC
MATERIALS, Dec. 16, 2014, KTH**

Sponsored by NSF-DMS-0920852 (Math. Bio.) and NSF-DMS-1413039 (Inverse Problems)



Poroelastic Material

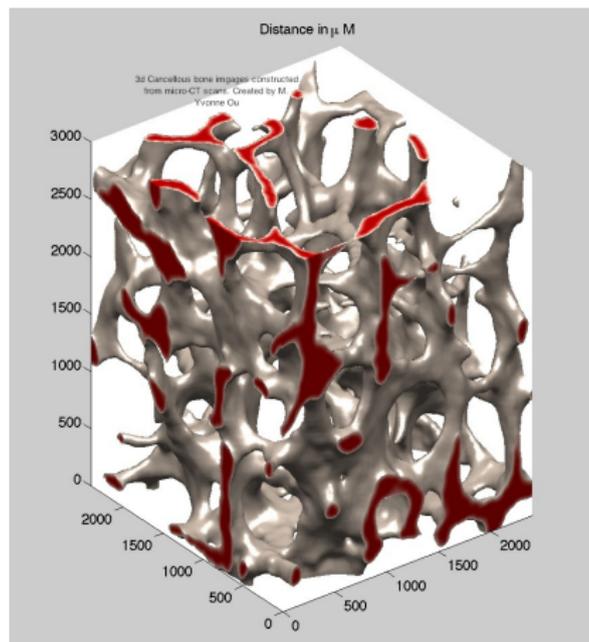


Figure: 3-D trabeculae reconstructed from micro-CT scans provided by L. Cardoso. [MYO-2011]. [Elastic frame with viscous pore fluid](#)

Outline

- Wave propagation in poroelastic composites and frequency-dependent mass coupling and dissipation (drag force).
- Stieltjes functions and the Integral Representation formula (IRF) for **dynamic permeability** with measure $d\lambda$
- Reconstruction of dynamic permeability from finite data sets
- IRF for **dynamic tortuosity** with measure $d\sigma$.
- Relations between the microstructure and tortuosity in terms of the moments of measure $d\sigma$

Biot's equations for waves in poroelastic materials

M. A. Biot, *Theory of propagation of elastic waves in a fluid-saturated porous solid*, J. Acoustical Society of America (JASA) (1956a (5510), 1956b (2412)).

Web of Science: We are sorry...The citation map you are trying to create contains more papers (nodes) than we can process.

Three pieces of physics

- Strain-stress relation
- Kinetic energy
- Energy dissipation

\mathbf{u} : displacement vector for **solid**

\mathbf{U} : displacement vector for **fluid**

* Stress tensor acting on the solid part: τ_{ij} , $i, j = 1, 2, 3$

* Stress tensor acting on the fluid part: $s\delta_{ij}$, $s = -\beta p$, β : **porosity**.

* Strain tensor in the solid: $e_{ij} = \frac{1}{2}(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i})$, $i, j = 1, 2, 3$

* Strain tensor in the fluid: $\epsilon = \nabla \cdot \mathbf{U}$

Coupled stress-strain relation $\tau_{ij} = 2N e_{ij} + \delta_{ij}(Ae + Q\epsilon)$

2nd-order formulation, low-frequency regime, isotropic case

Biot's Dynamic equations, 1956-a

$$N\nabla^2 \mathbf{u} + \nabla[(A + N)e + Q\epsilon] = \frac{\partial^2}{\partial t^2}(\rho_{11}\mathbf{u} + \rho_{12}\mathbf{U}) + \mathbf{b} \frac{\partial}{\partial t}(\mathbf{u} - \mathbf{U})$$

$$\nabla[Qe + R\epsilon] = \frac{\partial^2}{\partial t^2}(\rho_{12}\mathbf{u} + \rho_{22}\mathbf{U}) - \mathbf{b} \frac{\partial}{\partial t}(\mathbf{u} - \mathbf{U})$$

Dissipation coeff. $b := \frac{\beta^2 \eta}{\kappa}$ is constant \iff laminar flow in pore space

β : porosity, η : dynamic viscosity

κ : static permeability

First order formulation, low-freq.

$$\mathbf{v} := \dot{\mathbf{u}}, \quad \mathbf{q} := \beta(\dot{\mathbf{U}} - \dot{\mathbf{u}}), \quad p := -\frac{Qe + R\epsilon}{\beta} \text{ (pore pressure)}$$

- Stress-strain relations

$$\partial_t \tau_{xx} = c_{11}^u \partial_x v_x + c_{13}^u \partial_z v_z + \alpha_1 M (\partial_x q_x + \partial_z q_z)$$

$$\partial_t \tau_{zz} = c_{13}^u \partial_x v_x + c_{33}^u \partial_z v_z + \alpha_3 M (\partial_x q_x + \partial_z q_z)$$

$$\partial_t \tau_{xz} = c_{55}^u (\partial_z v_x + \partial_x v_z)$$

$$\partial_t p = -\alpha_1 M \partial_x v_x - \alpha_3 M \partial_z v_z - M (\partial_x q_x + \partial_z q_z)$$

- Equations of motion

$$\rho \partial_t v_i + \rho_f \partial_t q_i = [\nabla \cdot \tau]_i$$

(Darcy's law in x_i direction)

$$\rho_f \partial_t v_i + m_i \partial_t q_i + \left(\frac{\eta}{\kappa_i} \right) q_i = -\partial_{x_i} p$$

$m_i := \frac{\rho_f \alpha_{\infty i}}{\beta}$, $\alpha_{\infty i}$: inf-freq tortuosity in the i -th direction (difficult to measure)

Low-freq. Biot equations and ultrasound in cancellous bones

Problem: low-frequency Biot equations underestimate wave dissipation when compared with experiment.

Solution: High-frequency correction is made in frequency domain.

Fourier transform to frequency domain

$$\tilde{f}(\omega) := \frac{1}{\sqrt{2\pi}} \int_0^{\infty} v(t) e^{i\omega t} dt$$

Three approaches to high-freq. corrections

- Biot's correction to \tilde{b} : too restrictive in terms of pore geometry
- **JKD correction** introducing dynamic tortuosity and permeability function: a tunable parameter for geometry thus more flexible. However, the two JKD functions are the **simplest** form of the functions that satisfy the necessary causality properties. **There is no reason all pore geometries correspond to these specific function forms.**
- **Torquato-Avellaneda's integral representation formula (IRF)** for the dynamic permeability function: valid for all pore geometries **The most general correction.**

For the first order formulation, high-freq. correction is made to the dynamic tortuosity function.

Tortuosity and permeability, M. Matyka and Z. Koza, 2012

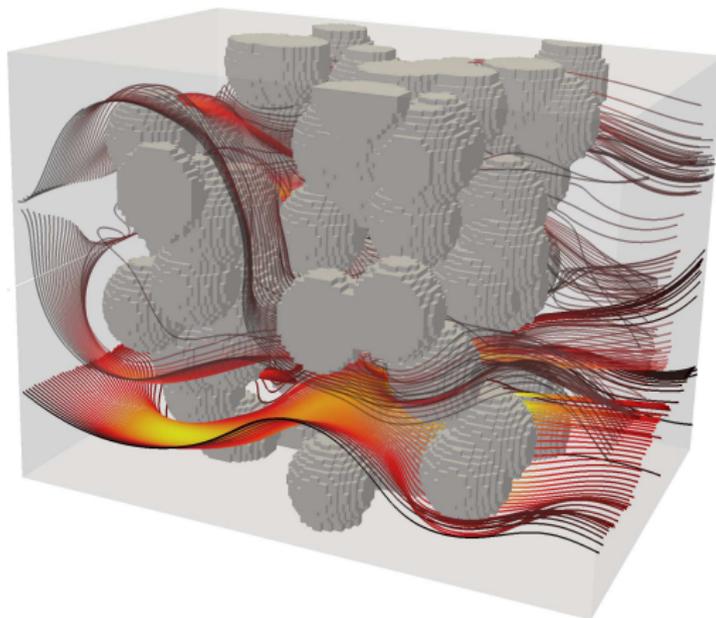


FIGURE 1. Streamlines in the fluid flow through three-dimensional random model of porous media at porosity $\phi = 0.6$ and tortuosity $T=1.15$.

Dynamic permeability and tortuosity

Dynamic permeability function $K(\omega)$

$$-i\omega\beta(\tilde{\mathbf{U}} - \tilde{\mathbf{u}}) = \frac{K(\omega)}{\eta}(-\nabla\tilde{p} + \rho_f\omega^2\tilde{\mathbf{u}}) \quad (1)$$

Dynamic tortuosity function $\alpha(\omega)$

$$\alpha(\omega)\rho_f(-i\omega)^2(\tilde{\mathbf{U}} - \tilde{\mathbf{u}}) = (-\nabla\tilde{p} + \rho_f\omega^2\tilde{\mathbf{u}}), \quad (2)$$

Example: (low-freq. Darcy's law)

$$\rho_f\partial_t v_x + \left(\frac{\rho_f\alpha_{\infty 1}}{\beta}\right)\partial_t q_x + \left(\frac{\eta}{\kappa_1}\right)q_x = -\partial_x p \iff \alpha_1(\omega) = \alpha_{\infty 1} + \frac{\eta\beta/(\kappa_1\rho_f)}{-i\omega}$$

Time-domain High-freq. Darcy's law has a memory term

$$(2) \text{ and } (1) \Rightarrow \alpha(\omega) = \frac{i\eta\beta}{\omega\rho_f}K^{-1}(\omega) \text{ for } \omega \neq 0$$

Note: $K(0) = \kappa = K_0$ =static permeability

JKD Model

In [Johnson-Koplik-Dashen-1987], 1077, by extending $\alpha(\omega)$ and $K(\omega)$ to complex ω -plane and using causality argument, the **simplest** forms are derived

$$\alpha_D(\omega) = \alpha_\infty \left(1 - \frac{\eta\phi}{i\omega\alpha_\infty\rho_f K_0} \sqrt{1 - i\frac{4\alpha_\infty^2 K_0^2 \rho_f \omega}{\eta\Lambda^2 \phi^2}} \right)$$

$$K_D(\omega) = K_0 / \left(\sqrt{1 - \frac{4i\alpha_\infty^2 K_0^2 \rho_f \omega}{\eta\Lambda^2 \phi^2}} - \frac{i\alpha_\infty K_0 \rho_f \omega}{\eta\phi} \right).$$

with the tunable geometry-dependent constant Λ , $\phi = \beta$ (porosity).
inf-freq. model as $\omega \rightarrow \infty$

$$\alpha(\omega) \rightarrow \alpha_\infty \left(1 + \sqrt{\frac{i\eta}{\rho_f \omega} \frac{2}{\Lambda}} \right), \quad K(\omega) \rightarrow \frac{i\eta\phi}{\alpha_\infty \rho_f \omega} \left(1 - \sqrt{\frac{i\eta}{\rho_f \omega} \frac{2}{\Lambda}} \right)$$

Models for ultrasound in cancellous bone

Hosakawa 2006, *Ultrasonics*

Table 3

Biot's parameters of bovine cancellous bone models with parallel and perpendicular trabecular orientations

	Parallel	Perpendicular
Young's modulus of solid bone E_s	22.0 GPa	22.0 GPa
Poisson's ratio of solid bone ν_s	0.32	0.23
Bulk modulus of bone marrow K_f	2.0 GPa	2.0 GPa
Density of solid bone ρ_s	1960 kg/m ³	1960 kg/m ³
Density of bone marrow ρ_f	930 kg/m ³	930 kg/m ³
Poisson's ratio of trabecular frame ν_b	0.32	0.23
Variable r	0.25	0.25
Viscosity of bone marrow η	1.5 P s	1.5 P s
Parameter n	1.46	2.14
Pore size a_0	1.35 mm	1.35 mm
Permeability k	3×10^{-8}	3×10^{-9}
Porosity β	0.83	0.83
Resistance coefficients γ_{ii} ($i = x, y$)	1.0×10^6	1.0×10^6
Resistance coefficient γ_{xy}	1.0×10^7	1.0×10^7

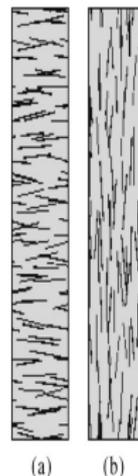


Fig. 2. Geometric drawings of bovine cancellous bone models in the viscoelastic FDTD method, with (a) parallel and (b) perpendicular trabecular orientations.

Fellah et. al. 2004, 2006, *JASA*: inf-freq. model.

Biot's model parameters of cancellous bone	M1	M2	M3
Thickness L (cm)	0.7	0.5	0.38
Density ρ_s (kg/m^3)	1960	1960	1960
Porosity ϕ	0.83	0.77	0.88
Tortuosity α_∞	1.05	1.01	1.02
Viscous characteristic length Λ (μm)	5	2.7	5
Bulk modulus of the elastic solid K_s (GPa)	20	20	26
Bulk modulus of bone skeletal frame K_b (GPa)	3.3	4	1.3
Shear modulus of the frame N (GPa)	2.6	1.7	0.35

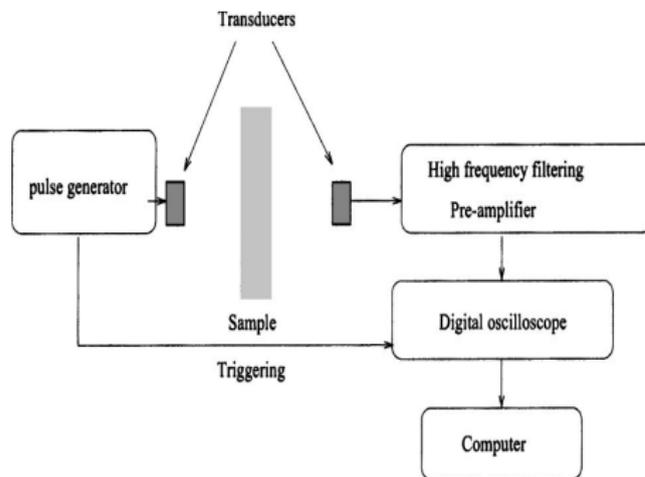


FIG. 15. Experimental setup for ultrasonic measurements.

In Fellah 2006, ϕ , α_∞ , Λ , K_b and N are sought after by solving the inverse problem.

Questions....

- How does the microstructural information of pore space **quantitatively** enter the permeability?
- Must all permeabilities have the simplest form as suggested in the JKD model? **No.**
- Given a set of permeability data at various frequencies, how well can we interpolate/extrapolate?
- How to numerically handle the memory terms in a time domain solver?

The key is in the mathematical structure of $K(\omega)$

Integral representation formula for $K(\omega)$

[Avellaneda-Torquato-91, 133]

For any pore geometry,

$$\frac{F}{\nu} K(\omega) = \int_0^{\Theta_1} \frac{\Theta dG(\Theta)}{1 - i\omega\Theta}, \quad \text{with } G(\Theta) = \frac{\sum_{\Theta_n \leq \Theta} b_n^2}{\sum_{n=1}^{\infty} b_n^2}$$

Note: dG is a probability measure (positive measure), which contains all geometric information of the pore space. $F = \frac{\alpha_\infty}{\phi}$, ν : kinetic viscosity. b_n depends on the spectrum of the Stokes equation in pore space and the initial applied flow direction e .

$$\Delta \Psi_n + \nabla Q_n = -\epsilon_n \Psi_n \text{ and } \nabla \cdot \Psi_n = 0 \text{ in } \mathcal{V}_1, \quad \Psi_n = 0 \text{ on } \partial\mathcal{V}_1,$$

$0 < \epsilon_1 \leq \epsilon_2 \leq \dots$ and $\epsilon_n \rightarrow \infty$ as $n \rightarrow \infty$. $\Theta_n := (\nu\epsilon_n)^{-1}$ are the viscous relaxation times and Θ_1 referred to as the principal viscous relaxation time. The eigenfunctions are orthonormal in the sense

$$\frac{1}{|\mathcal{V}_1|} \int_{\mathcal{V}_1} \Psi_m(\mathbf{x}) \cdot \Psi_n(\mathbf{x}) d\mathbf{x} = \delta_{mn} \text{ (Kronecker delta)}, \quad b_n = \frac{1}{|\mathcal{V}_1|} \int_{\mathcal{V}_1} e \cdot \Psi_n(\mathbf{x}) d\mathbf{x} \quad (3)$$

New Variables

ω	$s := -i\omega$	$\xi := -\frac{1}{s}$
$K(\omega)$	$P(s) := \left(\frac{F}{\nu}\right) K(is)$	$R(\xi) := -sP(s)$

$$P(s) := \left(\frac{F}{\nu}\right) K(is) = \int_0^{\Theta_1} \frac{\Theta dG(\Theta)}{1 + s\Theta} =: \int_0^{\Theta_1} \frac{d\lambda(\Theta)}{1 + s\Theta}$$

$$R(\xi) := -s \left(\frac{F}{\nu}\right) K(is) = -sP(s) = \int_0^{\Theta_1} \frac{\Theta dG(\Theta)}{\xi - \Theta}$$

Note: JKD model satisfies this IRF as well (can be shown using Herglotz function theory [MYO-2014])

Stieltjes function

Definition (Stieltjes function)

A Stieltjes function $f(z)$ for z on the extended complex plane has the following form

$$f(z) = \int_a^b \frac{d\mu(t)}{z - t} \quad (4)$$

where a, b are extended real numbers and $\mu(t)$ is a bounded, non-decreasing real function.

$$R(\xi) := -s \left(\frac{F}{\nu} \right) K(is) = -sP(s) = \int_0^{\Theta_1} \frac{\Theta dG(\Theta)}{\xi - \Theta}$$

is a Stieltjes function. Note: $\xi = -\frac{i}{\omega}$.

Multipoint Padé approximates of Stieltjes functions

Theorem ([Gelfgren-1978])

Let f be a Stieltjes function. Let $P_{n-1}(z)$ and $Q_n(z)$ be polynomials of degree at most $n-1$ and n , respectively, satisfying the relations ($k_1 + k_2 + k_3 = 2n$)

$$\begin{cases} f(z)Q_n(z) - P_{n-1}(z) = A(z) \prod_{j=1}^{k_1} (z - x_j) \prod_{j=1}^{k_2} (z - z_j)(z - \bar{z}_j) \\ f(z)Q_n(z) - P_{n-1}(z) = B(z)z^{n-k_3-1} \end{cases}$$

where $A(z)$, $B(z)$ are analytic in $\mathbb{C} \setminus [a, b]$, $B(z)$ bounded at ∞ , $x_1, \dots, x_{k_1} \in \mathbb{R} \setminus [a, b]$, $z_1, \dots, z_{k_2} \in \mathbb{C} \setminus \mathbb{R}$. Then

$[n-1/n]_f(z) := \frac{P_{n-1}(z)}{Q_n(z)} = \int_a^b \frac{d\beta(t)}{z-t}$ for some bounded, non-decreasing function $\beta(t)$.

$[n-1/n]_f(z) = \sum_{k=1}^n \frac{r_k}{z-p_k}$, with $r_k > 0$ and $p_k \in (a, b)$ for $k = 1, \dots, n$.

Reconstruction from finite data set $K(\omega_j)$

Given $R(\xi_j)$, $j = 1, \dots, M$, to construct $[M - 1/M]_R(\xi)$. **Data at conjugate points can be obtained through symmetry**

$$R(\bar{\xi}_j) = \int_0^{\Theta_1} \frac{d\lambda(\theta)}{\bar{\xi}_j - \theta} = \overline{\int_0^{\Theta_1} \frac{d\lambda(\theta)}{\xi_j - \theta}} = \overline{R(\xi_j)}$$

$$[M - 1/M]_R(\xi) = \sum_{j=1}^M \frac{r_j}{\xi - p_j}, \quad r_j > 0, 0 < p_j < \Theta_1$$

$$K(\omega) \approx \frac{\nu}{F} \sum_{j=1}^M \frac{r_j}{1 - i\omega p_j}, \quad \mathcal{F}^{-1}[K](t) \approx \left(\frac{\nu}{F}\right) \sum_{j=1}^M \left(\frac{r_j}{p_j}\right) e^{-\frac{t}{p_j}}$$

Moments: $\mu_j(d\lambda) := \int_a^b \theta^j d\lambda(\theta) \approx \sum_k p_k^j r_j$

Note: $\mu_0(d\lambda) = \frac{FK_0}{\nu} = \frac{\alpha_\infty K_0}{\nu\phi} \Rightarrow \alpha_\infty$ can be computed from permeability data

IRF for Dynamic Tortuosity

Theorem (MYO-2014)

The dynamic tortuosity $\alpha(\omega) = \frac{i\eta\phi}{\omega\rho_f} K^{-1}$ has the following integral representation formula for ω such that $-\frac{i}{\omega} \in \mathbb{C} \setminus [0, \Theta_1]$

$$\alpha(\omega) = a \left(\frac{i}{\omega} \right) + \int_0^{\Theta_1} \frac{d\sigma(\Theta)}{1 - i\omega\Theta}$$

for some positive measure $d\sigma$, with $a = \frac{\eta\phi}{\rho_f K_0}$. $\mu_0(d\sigma) = \frac{\alpha_\infty}{\mu_0^2(d\lambda)} \mu_1(d\lambda)$

- (1) $\alpha(\omega) \rightarrow \alpha_\infty$ as $\omega \rightarrow \infty$, $d\sigma$ has a Dirac mass at $\Theta = 0$ with strength α_∞ .
- (2) $h(\xi) := a - \frac{\alpha_\infty}{\xi R(\xi)} = a + i\omega\alpha(\omega)$ is a Stieltjes function and $h(\bar{\xi}) = \overline{h(\xi)}$
- (3) $\alpha(\omega) \approx \frac{a}{-i\omega} + \sum_{j=1}^M \frac{r_j}{-i\omega - p_j}$, $r_j > 0$, $p_j \in (-\infty, -\frac{1}{\Theta_1})$
- (4) $\mathcal{F}^{-1}[\alpha](t) = a\delta(t) + \sum_{j=1}^M r_j e^{p_j t}$ Provides an efficient way to handle the memory term and settles a debate in poroelasticity

Moments and effective properties [MYO-2014]

$$\mu_0(d\lambda) = \frac{\alpha_\infty K_0}{\phi\nu}$$

$$\alpha_\infty = \frac{\mu_0(d\sigma)\mu_0^2(d\lambda)}{\mu_1(d\lambda)}$$

Specializing the result to JKD model, the microstructure-dependent parameter Λ (a weighted average of volume-to-surface ratio of the dynamically connected pore space) satisfies

$$\Lambda = \sqrt{\frac{2K_0\alpha_\infty^2}{\phi[\mu_0(d\sigma^D) - \alpha_\infty]}} = \sqrt{\frac{2K_0\alpha_\infty}{\phi\left[\frac{\mu_1(d\lambda^D)}{\mu_0^2(d\lambda^D)} - 1\right]}}$$

Note: Empirical formula in literature $\Lambda \approx \sqrt{\frac{2\alpha_\infty K_0}{\phi/4}}$!

Formulation for Reconstruction

$$P(s_j) = \int_0^{\Theta_1} \frac{\Theta dG(\Theta)}{1+s_j\Theta} \approx [M - 1/M]P(s) := \frac{a_0 + a_1 s_j + \dots + a_{M-1} s_j^{M-1}}{1 + b_1 s_j + \dots + b_M s_j^M}, \quad j = 1 \sim$$

$2M$ Given the data $(s_j, P(s_j))$, $s_j \neq 0$, $j = 1, \dots, M$, let $s_{j+M} = -s_j$, $P(s_{j+M}) = \overline{P(s_j)}$ and $P_m := P(s_m)$, $m = 1, \dots, 2M$. Solve $A\mathbf{x} = \mathbf{d}$

$$A = \begin{pmatrix} 1 & s_1 & s_1^2 & \dots & s_1^{M-1} & -P_1 s_1 & -P_1 s_1^2 & -P_1 s_1^3 & \dots & -P_1 s_1^M \\ 1 & s_2 & s_2^2 & \dots & s_2^{M-1} & -P_2 s_2 & -P_2 s_2^2 & -P_2 s_2^3 & \dots & -P_2 s_2^M \\ \vdots & \vdots \\ 1 & s_{2M} & s_{2M}^2 & \dots & s_{2M}^{M-1} & -P_{2M} s_{2M} & -P_{2M} s_{2M}^2 & -P_{2M} s_{2M}^3 & \dots & -P_{2M} s_{2M}^M \end{pmatrix}$$

$$= A_r + i A_i := \text{Real}(A) + i \text{Imag}(A)$$

$$\mathbf{x} = (a_0, a_1, \dots, a_{M-1}, b_1, b_2, \dots, b_M)$$

$$\mathbf{d} = (P_1 \ P_2 \ \dots \ P_{2M})^t$$

$$[M - 1/M]P(s) := \frac{a_0 + a_1 s + \dots + a_{M-1} s^{M-1}}{1 + b_1 s + \dots + b_M s^M} = \sum_{j=1}^M \frac{r_j}{s - p_j} \Rightarrow \mathbf{x} \in \mathbb{R}^{2M}$$

Algorithm [MYO-2014]

- ① Rescale each column \mathbf{a}_j , $j = 1, \dots, 2M$ of A by $C := \text{diag}\{\|\mathbf{a}_j\|_2^{-1}\}_{j=1}^{2M}$.
Let

$$A\mathbf{x} = (AC)(C^{-1}\mathbf{x}) =: B\mathbf{y}$$

- ② Let

$$\hat{B} := \begin{pmatrix} B_r \\ B_i \end{pmatrix}, \quad \hat{\mathbf{d}} := \begin{pmatrix} \mathbf{d}_r \\ \mathbf{d}_i \end{pmatrix}.$$

Solve the overdetermined system $\hat{B}\mathbf{y} = \hat{\mathbf{d}}$ in least square sense with Tikhonov regularization

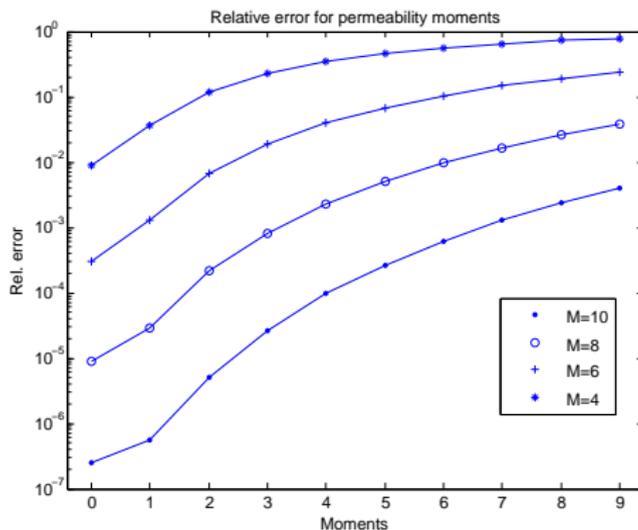
$$\min_{\mathbf{y}} \|\hat{B}\mathbf{y} - \hat{\mathbf{d}}\|_2^2 + \gamma^2 \|\mathbf{y}\|_2^2$$

L -curve method is used for choosing the regularization parameter γ

- ③ Rescale $\mathbf{x} = C\mathbf{y}$.
- ④ Apply partial fraction decomposition to the resulting $2M$ Padé approximant. Retain the poles in $(-\infty, -\frac{1}{C_1}] \cup \{-\frac{1}{\xi_p}\}$ with positive weights.

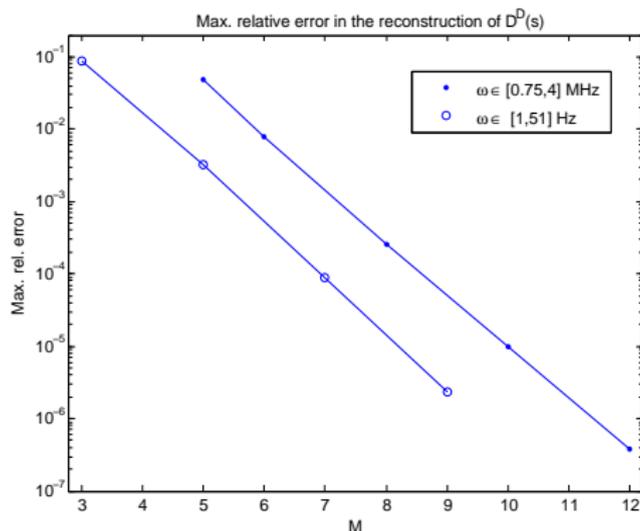
JKD permeability moments from data $K(\omega_j)$ [MYO-2014]

Figure: Relative error of $\mu_1(d\lambda^D), \dots, \mu_9(d\lambda^D)$ approximated from M data points in $[1, 51]$ Hz



Reconstructed JKD Tortuosity from data $K(\omega_j)$ [MYO-2014]

Figure: . Max. Rel. Error for $\omega \in [1, 51]$ Hz and $\omega \in [0.75, 4]$ MHz. Linear convergence!



On reconstruction of dynamic permeability and tortuosity from data at distinct frequencies, *Inverse Problems* 30(9) 095002, 2014